# Head-on Collision Experiment of Pulse-like Nonlinear Surface Acoustic Waves

パルス状非線形弾性表面波の正面衝突実験

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## 1. Introduction

Nonlinear surface acoustic wave (NLSAW) generated on the piezoelectric crystal has attracted much attention in application to signal processing devices such as convolver, parametric amplifier, and memory correlator. [1-3]

Recently, the dynamic characteristics of NLSAW have been studied using a complex aproach by authors. [4-9] In this approach, dynamic behavior of the fundamental (FUN) and higher-order harmonics (H-HMCs) in NLSAW and their interaction during propagation were investigated by optical diffraction technique. In the collision experiment of two 2<sup>nd</sup> HMCs, the diffracted light intensities were preserved before and after collision. [7] To understand the dynamic characteristics of the 2<sup>nd</sup> HMC with dispersive medium, the collision experiment of two  $2^{nd}$  HMCs on Au film with various thicknesses was carried out. [8] As an application using dymanic characteristics of the 2<sup>nd</sup> HMC, authors have already proposed a novel photonic signal control device. [9] However, further experimental investigation of nonlinear behavior has been needed in order to comprehend the dynamics of NLSAW. In this study, we investigated the variation of the phase velocity and the full-width at half maximum (FWHM) value of a pulse-like NLSAW propagating on LiNbO<sub>3</sub> with Au film through a head-on collision experiment.

## 2. Experimental procedure

Figure 1 shows an illustration of an experimental sample. 128°-rotated Y-cut X-propagating LiNbO<sub>3</sub> single crystal was used as substrate. The 20 pair interdigital transducer (IDT) was formed on the substrate, in which their center frequency and insertion loss were 48.7 MHz and 6.85 dB, respectively. A space between IDT(I) and (II) was fixed at 20 mm. The RF burst signal was simultaneously introduced to two IDTs through RF power amplifier from synthesized RF signal generator, in which its duration time and duty cycle were 250 ns and 0.05 %, respectively. 50-nm-thick Au film as dispersive material was

IDT(I) He-Ne Laser IDT(I) IDT(I) IDT(I) IDT(I) IDT(I) IDT(I) IDT(I)

Fig.1 I Illustration of a fabricated SAW device and an optical diffraction method for detecting generated pulse-like NLSAW.

deposited on the center of the sample surface by vacuum evaporation, in which its area was 10x10 Each pulse-like NLSAW propagating  $\mathrm{mm}^2$ . toward the center of the device from IDT(I) and (II) is head-on collision at the center of Au film. The dynamic characteristics of the FUN and the 2<sup>nd</sup> HMC were measured using optical diffraction method. A He-Ne laser beam with 300 µm in diameter was used for probing the NLSAW. The probing beam was irradiated on the sample surface, and the diffracted light was detected by photomultiplier tube. The waveform shown in Fig.2 is observed as the envelope profile of NLSAW using digital oscilloscope.



Fig.2 Typical envelope profiles of FUN and 2<sup>nd</sup> HMC

## 3. Results and discussion

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Figure 3 shows the propagation of the pulse-like



Fig.3 Intensity plot of (a) FUN and (b)  $2^{nd}$  HMC showing propagation of pulse-like waves from IDT(I) and (II). Collision point is indicated by arrow.

FUN and the  $2^{nd}$  HMC, in which the input power densities of IDT(I) and (II) are 123 and 99 W/mm<sup>2</sup>, respectively. The diffracted light intensity was plotted as function of the propagation time and distance from the collision point. As can be seen from these figures, the diffracted light intensity of  $2^{nd}$  HMC at the collision point was larger than the algebraic sum of two pulses before and after collision, whereas its value of FUN was smaller. The diffracted light intensity variation at the collision point must be able to explain the enhancement of coupling efficiency between FUN and  $2^{nd}$  HMC and energy conversion from FUN to  $2^{nd}$  HMC.

From the experimental results shown in Fig.3, the phase velocity and the FWHM value can be estimated. Figure 4 shows the phase velocity and the FWHM value of the FUN and the 2<sup>nd</sup> HMC from the IDT(I) and (II) as a function of propagation distance from the collision point. It seems from Fig. 4(a) and (b) that the velocity perturbation is increased by a head-on collision. These results may suggest the interaction between SAW pulses by nonlinearlity. The FWHM values before and after collision were almost the same as shown in Fig. 4(c) and (d). However, the FWHM value at the collision point was quit different between the FUN and the  $2^{nd}$  HMC. The pulse width of the FUN at collision was wider than that before and after collision, whereas its value of the  $2^{nd}$  HMC was not varied by collision. It seems from these results that a pulse-like  $2^{nd}$  HMC is



Fig.4 Phase velocity variation of (a) the FUN and (b) the  $2^{nd}$  HMC from IDT(I) and (II) as a function of propagation distance from collision point. (c) and (d) show FWHM variations of the FUN and  $2^{nd}$  HMC, respectively. Collision point is indicated by arrow.

stable with respect to collision.

#### 4. Conclusion

The dynamic behavior of the phase velocity and the FWHM value of the pulse-like NLSAW was investigated by a head-on collision experiment. The phase velocity at the collision point was perturbed and then the FWHM value of the FUN slightly increased at the collision point, whereas its value of the  $2^{nd}$  HMC was unchanged by collision.

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